

### Maxwell Distribution of Particle Velocities

Given a set of particles,  $N$ , with a temperature,  $T$ , then according to the Boltzmann distribution, the number of particles,  $N_i$ , with a kinetic energy,  $E_i$ , is

$$\frac{N_i}{N} = \frac{g_i \exp\left(-\frac{E_i}{k_B T}\right)}{\sum_j g_j \exp\left(-\frac{E_j}{k_B T}\right)} \quad (1)$$

where  $g_i$  is the degeneracy, that is, the number of states having the same energy,  $E_i$ . The kinetic energy of a particle is

$$\frac{1}{2}mv^2 = \frac{1}{2}m(v_x^2 + v_y^2 + v_z^2) \quad (2)$$

In considering a large number of particles, we can generalize the degeneracy concept as follows. In 3-dimensional velocity space ( $v_x$ ,  $v_y$  and  $v_z$ ), the velocity vectors that correspond to a given speed,  $|v|$ , lie on the surface of a sphere with radius  $|v|$ . The larger  $|v|$  is, the bigger the sphere is, and the more possible velocity vectors there are. So the number of possible velocity vectors for a given speed scales with the surface area of a sphere of radius  $|v|$ . Therefore the probability density of the speed is

$$f(|\vec{v}|) = c \, 4\pi v^2 \exp\left(-\frac{m(v_x^2 + v_y^2 + v_z^2)}{2k_B T}\right) = c \, 4\pi v^2 \exp\left(-\frac{mv^2}{2k_B T}\right) \quad (3)$$

where  $\int f(|v|)dv = 1$ , that is the integral over all probability density must integrate to 1. We use this to determine the constant,  $c$ , in (3) so that (3) is properly normalized.

$$\int_0^{\infty} f(|\vec{v}|)dv = 1 = \int_0^{\infty} c \, 4\pi v^2 \exp\left(-\frac{mv^2}{2k_B T}\right)dv = c \, 4\pi \int_0^{\infty} v^2 \exp\left(-\frac{mv^2}{2k_B T}\right)dv \quad (4)$$

We use the definite integral

$$\int_0^{\infty} x^2 e^{-ax^2} dx = \frac{1}{4} \sqrt{\frac{\pi}{a^3}} \quad (5)$$

with  $a = m/(2k_B T)$  to get

$$\int_0^{\infty} f(|\vec{v}|)dv = 1 = c \, 4\pi \int_0^{\infty} v^2 \exp(-av^2)dv = \frac{c \, 4\pi}{4} \sqrt{\frac{\pi}{a^3}} = c \left(\frac{2\pi k_B T}{m}\right)^{3/2} \quad (6)$$

Therefore the normalization constant,  $c$ , is given by

$$c = \left(\frac{m}{2\pi k_B T}\right)^{3/2} \quad (7)$$

and the probability density of the particle thermal or kinetic speed is

$$f(|\vec{v}|) = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi v^2 \exp\left(-\frac{m(v_x^2 + v_y^2 + v_z^2)}{2k_B T}\right) \quad (8)$$

The *expected or mean value of*  $|\mathbf{v}|$  is therefore

$$\overline{|\mathbf{v}|} = \int_0^\infty v f(|\vec{v}|) dv = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi \int_0^\infty v^3 \exp\left(-\frac{mv^2}{2k_B T}\right) dv \quad (9)$$

Combining (9) with the definite integral

$$\int_0^\infty x^3 e^{-ax^2} dx = \frac{1}{2a^2} \quad (10)$$

and defining  $a = m/(2 k_B T)$  yields the mean of the kinetic speed

$$\overline{|\mathbf{v}|} = \int_0^\infty v f(|\vec{v}|) dv = \left(\frac{m}{2\pi k_B T}\right)^{3/2} \frac{4\pi}{2} \left(\frac{2k_B T}{m}\right)^2 = \left(\frac{8k_B T}{\pi m}\right)^{1/2} \quad (11)$$

The *most probable value* of  $|\mathbf{v}|$  is when  $f(|\mathbf{v}|)$  reaches a maximum where  $d\{f(|\mathbf{v}|)\}/d|\mathbf{v}| = 0$

$$\frac{df(|\vec{v}|)}{dv} = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi \left\{ 2v \exp\left(-\frac{mv^2}{2k_B T}\right) + v^2 \exp\left(-\frac{mv^2}{2k_B T}\right) \left[-\frac{mv}{k_B T}\right] \right\} = 0 \quad (12)$$

$$2v = v^2 \left[ \frac{mv}{k_B T} \right] \quad (13)$$

$$v = \sqrt{\frac{2k_B T}{m}} \quad (14)$$

The square root of the mean square or *rms velocity* is

$$\overline{v^2} = \int_0^\infty v^2 f(|\vec{v}|) dv = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi \int_0^\infty v^4 \exp\left(-\frac{mv^2}{2k_B T}\right) dv \quad (15)$$

Combining (15) with the definite integral

$$\int_0^\infty x^4 e^{-ax^2} dx = \frac{3}{2^3} \frac{\sqrt{\pi}}{a^{5/2}} \quad (16)$$

with  $a = m/(2 k_B T)$  yields

$$\overline{v^2} = \int_0^\infty v^2 f(|\vec{v}|) dv = \left(\frac{m}{2\pi k_B T}\right)^{3/2} 4\pi \frac{3\sqrt{\pi}}{8} \left(\frac{2k_B T}{m}\right)^{5/2} = \frac{3k_B T}{m} \quad (17)$$

$$v_{rms} = \left(\overline{v^2}\right)^{1/2} = \sqrt{\frac{3k_B T}{m}} \quad (18)$$